

# Fullqubit alchemist: *Quantum algorithm for alchemical free energy calculations*

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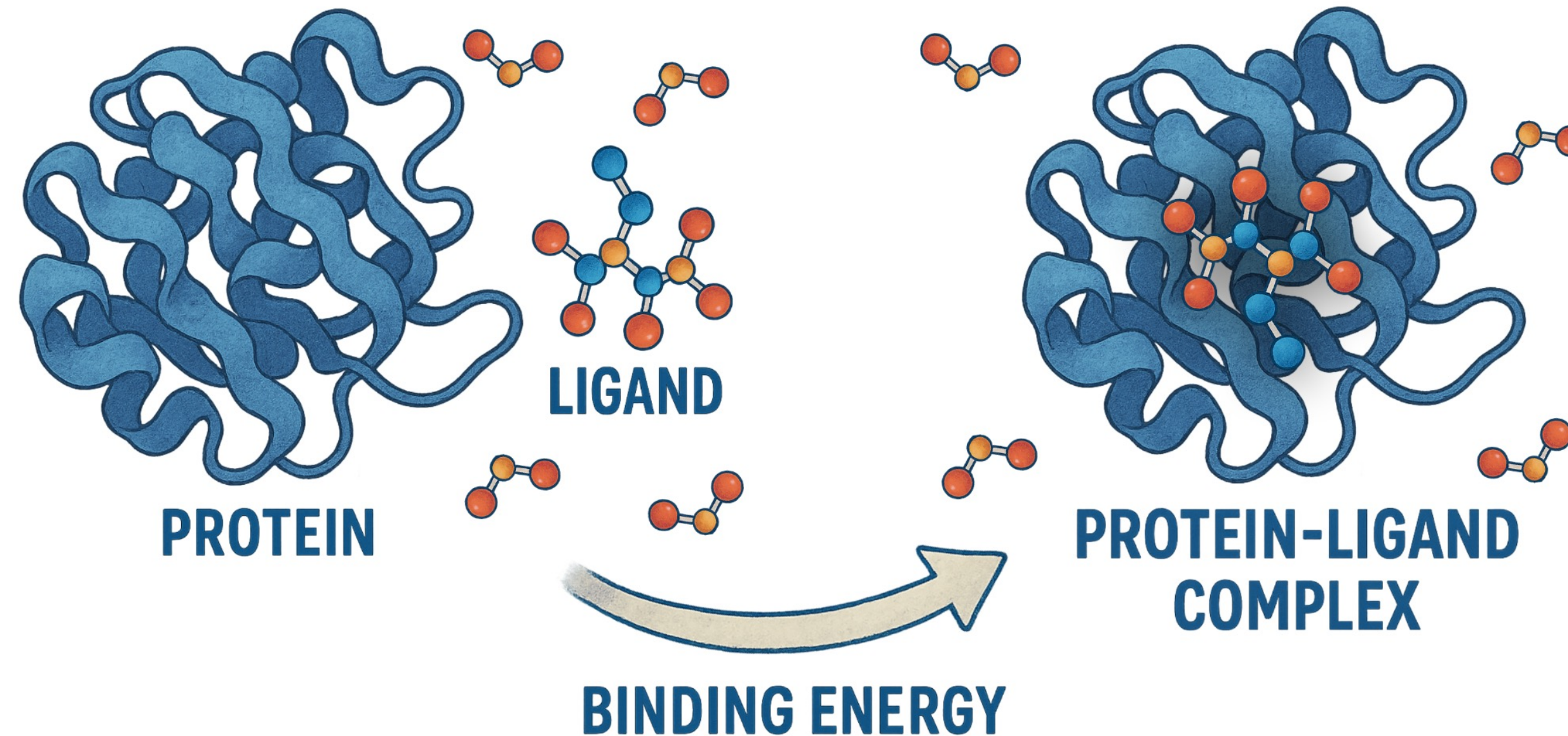


QUANTUM  
MOTION

# Drug design

## *Motivation and applications*

- Binding free energy is used to rank drug candidates.
- Free energy is computed via thermal averages over many many many different molecular configurations.



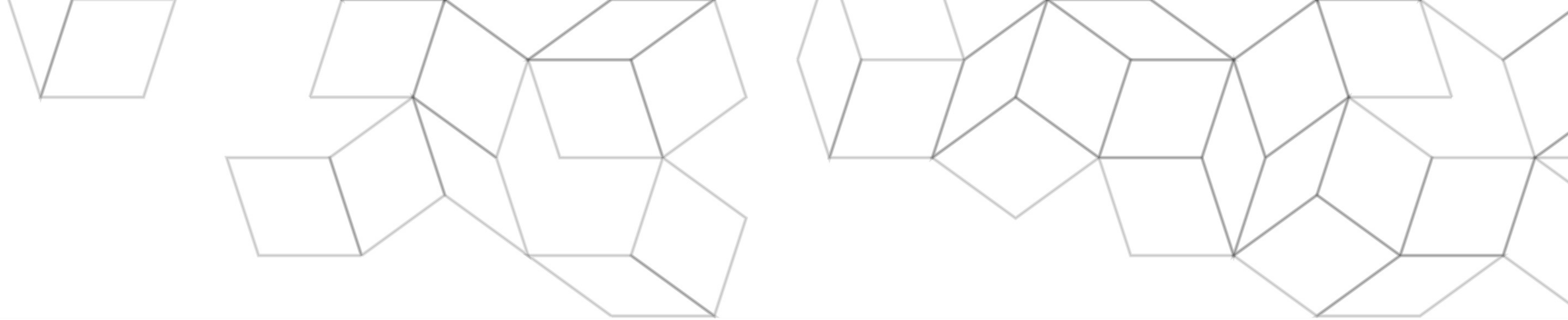
# Talk questions and outline

*PART I: Molecular dynamics on a quantum computer*

How do we generate the various molecular configurations on the quantum computer?

*PART II: Alchemical free energy calculations*

How do we calculate the relative free energy between two different systems?

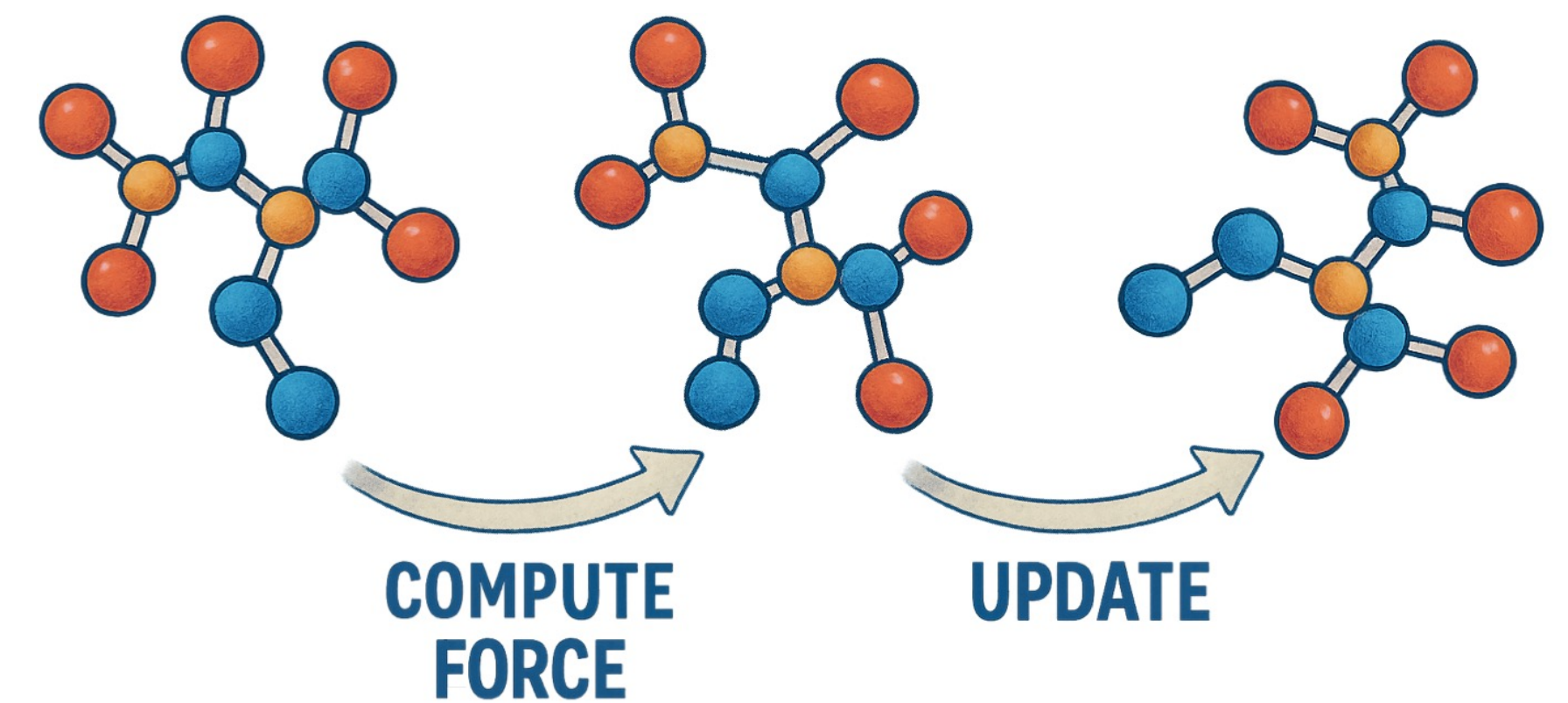


# **Part I: Molecular Dynamics on a Quantum Computer**

# Molecular dynamics

*In practice for classical computers*

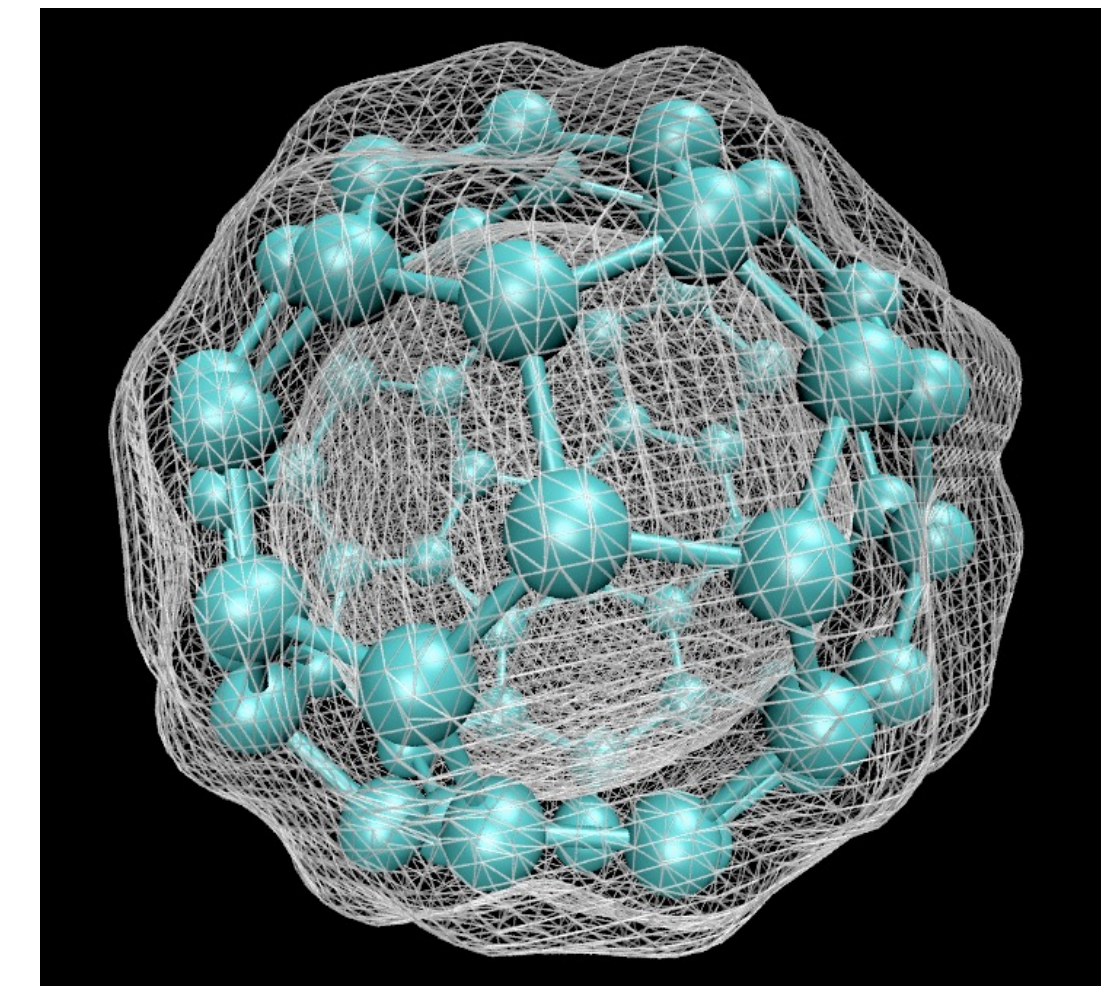
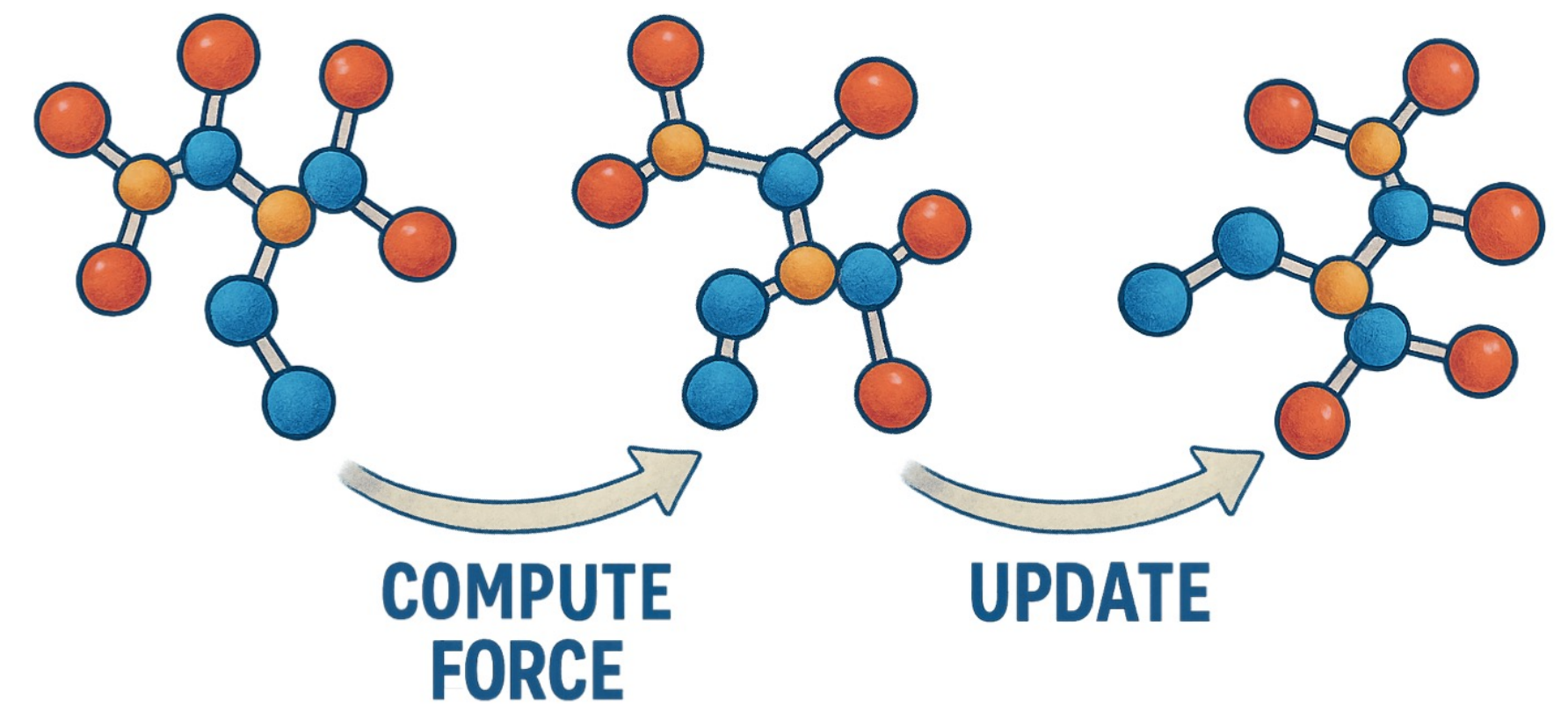
- Propagate particles with Hamiltonian dynamics
- Time average of evolution trajectory to replace ensemble average
  - Lots and lots of compute time!
- Used to compute thermal averages such as free energy



# Hybrid quantum semi-classical modeling

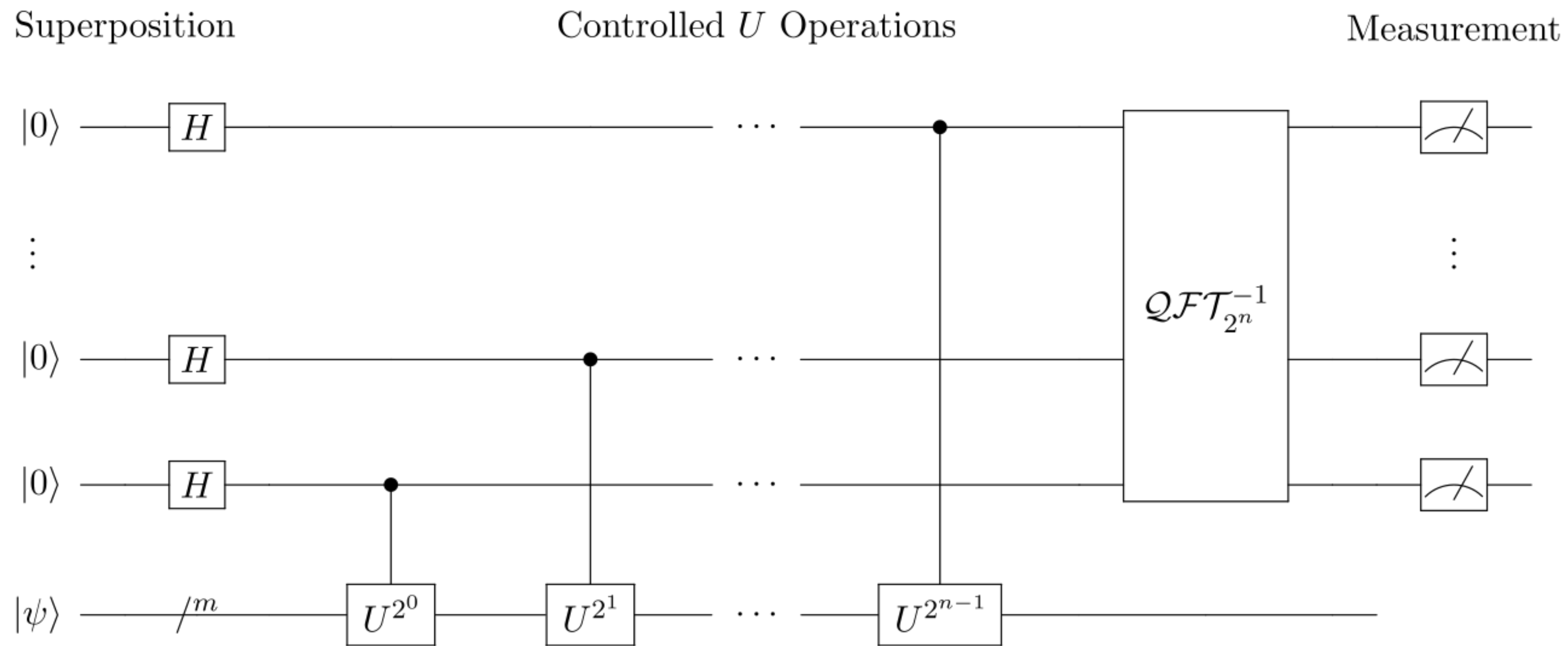
*Not all particles are classical...*

- While nuclei are heavy enough to be considered classical, electrons whizzing around definitely are not...
- Born-Oppenheimer Approximation:
  - The wave function of the *nuclei* and *electrons* are treated separately.
  - Treat nuclei *classically* → propagate by molecular dynamics
  - Treat electrons *quantumly* → electronic structure calculations (eg DFT, CCSD) of the ground state



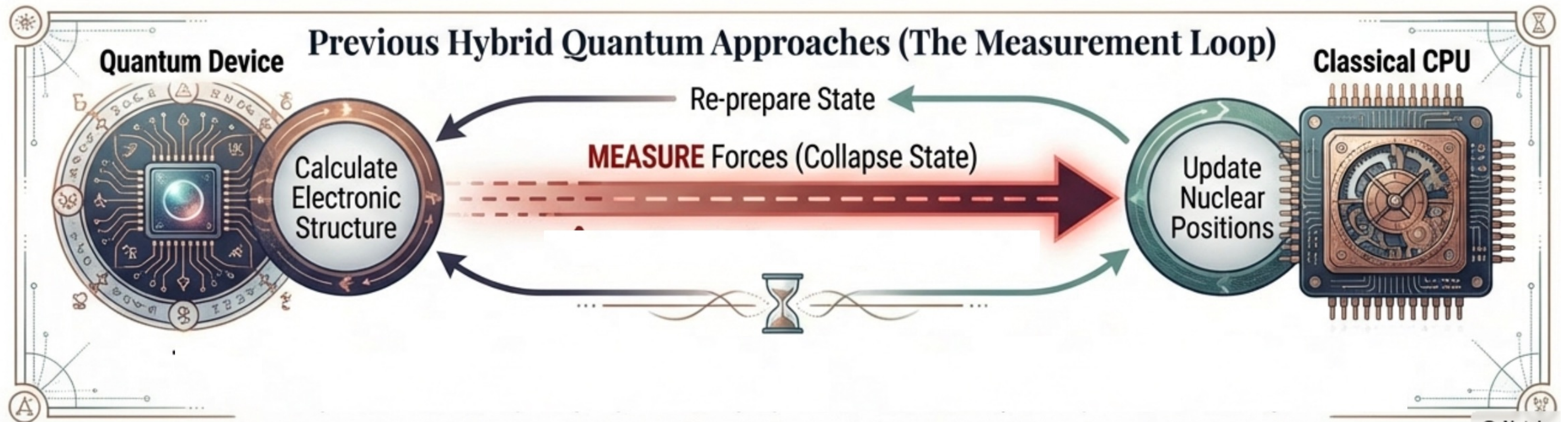
# Calculating ground state properties...

*Seem familiar?*



- Offload ground state energy calculations to quantum computer
- Feed ground state energy calculations to MD step

# Case closed?



Is there a coherent approach that doesn't require hybrid loops?  
Can we implement molecular dynamics on a quantum computer?

## Improved Precision Scaling for Simulating Coupled Quantum-Classical Dynamics

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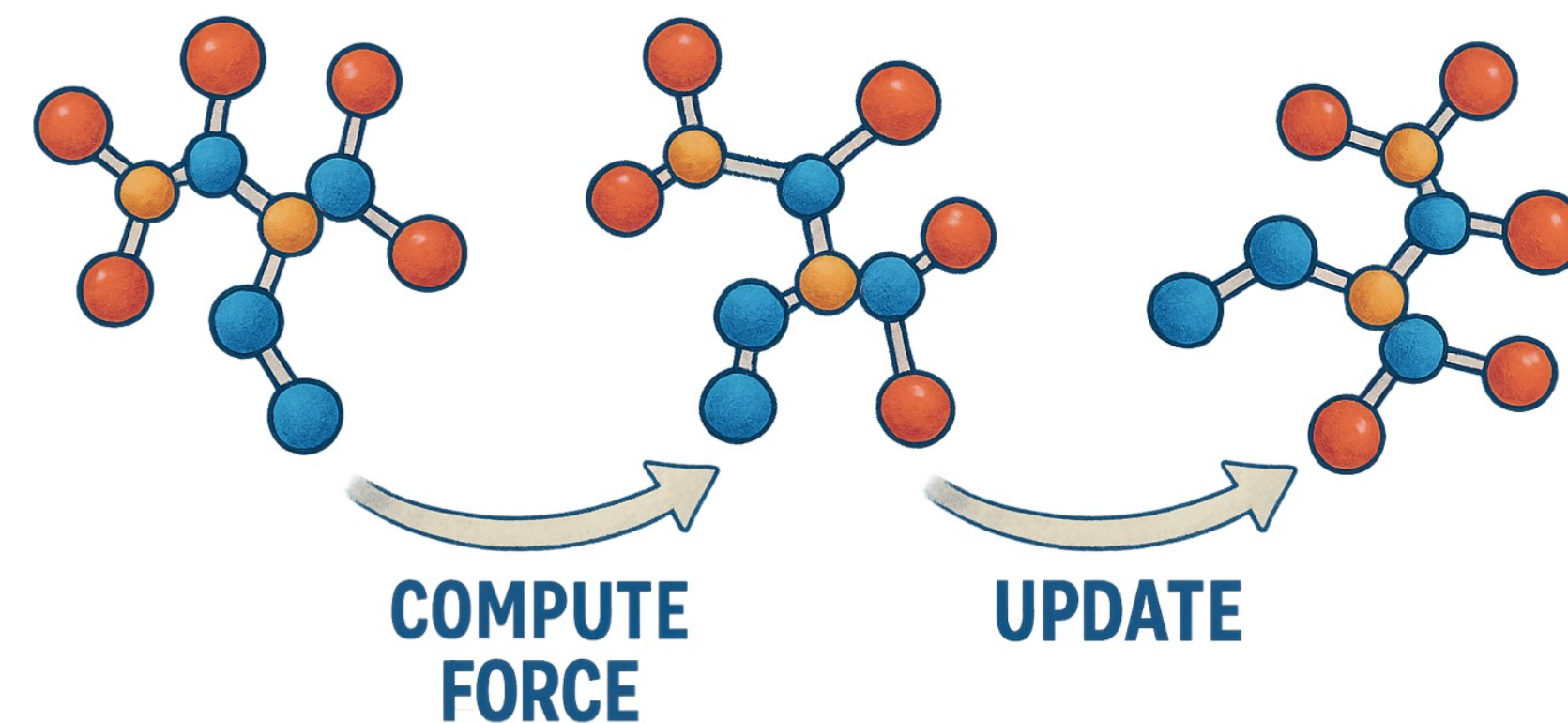


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# Molecular dynamics on a quantum computer

## *What works and what doesn't*

- Large Hilbert space with small number of qubits
- Can simulate/discretize the entire phase-space
- A quantum state can be used to encode an entire probability distribution
- Time-based averages may be hard to implement



# Koopman von Neumann mechanics

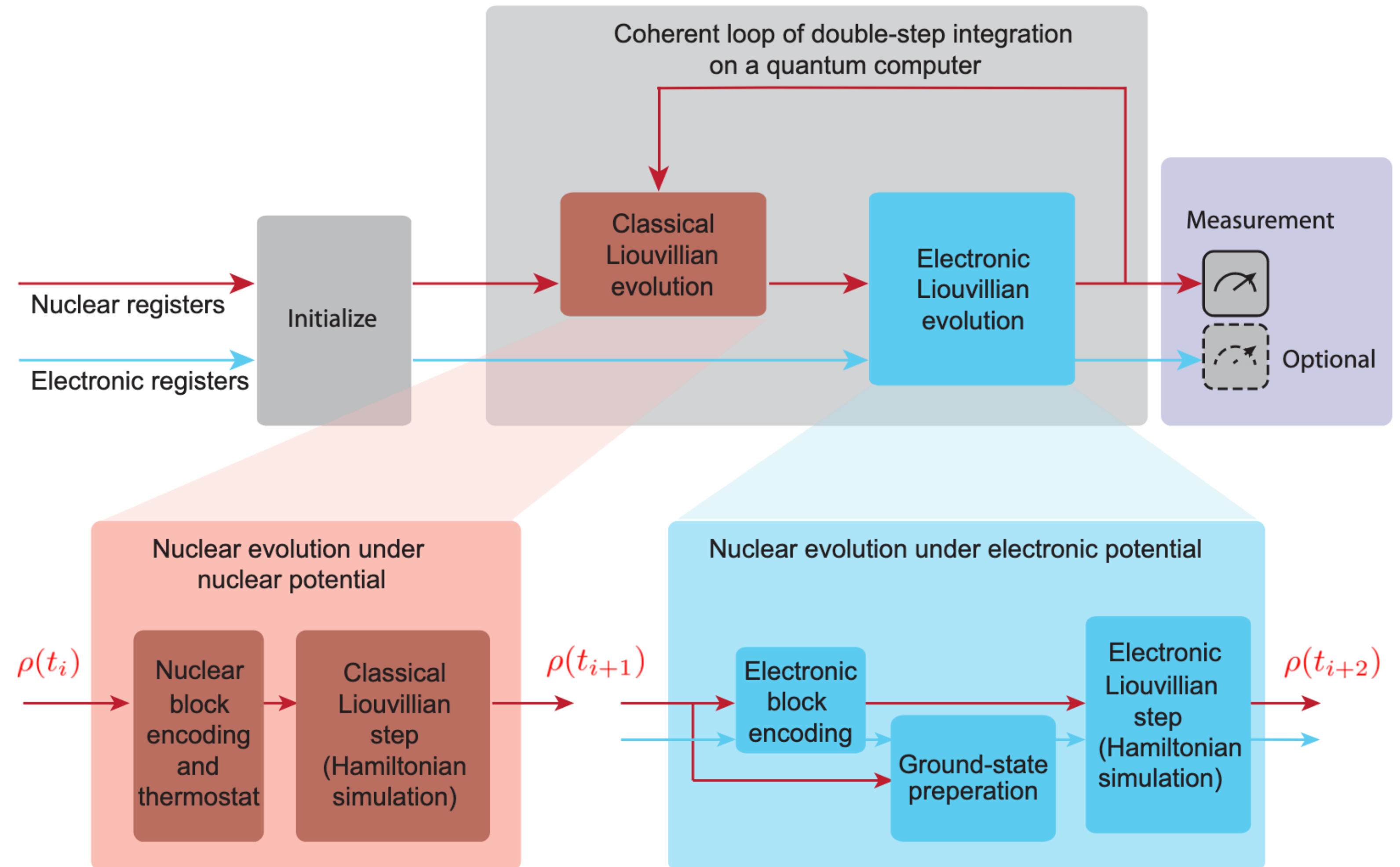
*Under the Liouvillian picture*

- Describes classical dynamics over a classical probability distribution of a statistical ensemble
- Classical wavefunction  $\rho = \psi_{KvN}^*(p, q, t)\psi_{KvN}(p, q, t)$  that encodes distribution
- Evolution of wavefunction by Liouvillian  $L = -i \left( \frac{\partial H}{\partial p} \frac{\partial}{\partial x} - \frac{\partial H}{\partial x} \frac{\partial}{\partial p} \right)$
- Evolution dynamics governed by  $i \frac{\partial}{\partial t} \psi_{KvN} = L\psi_{KvN}$

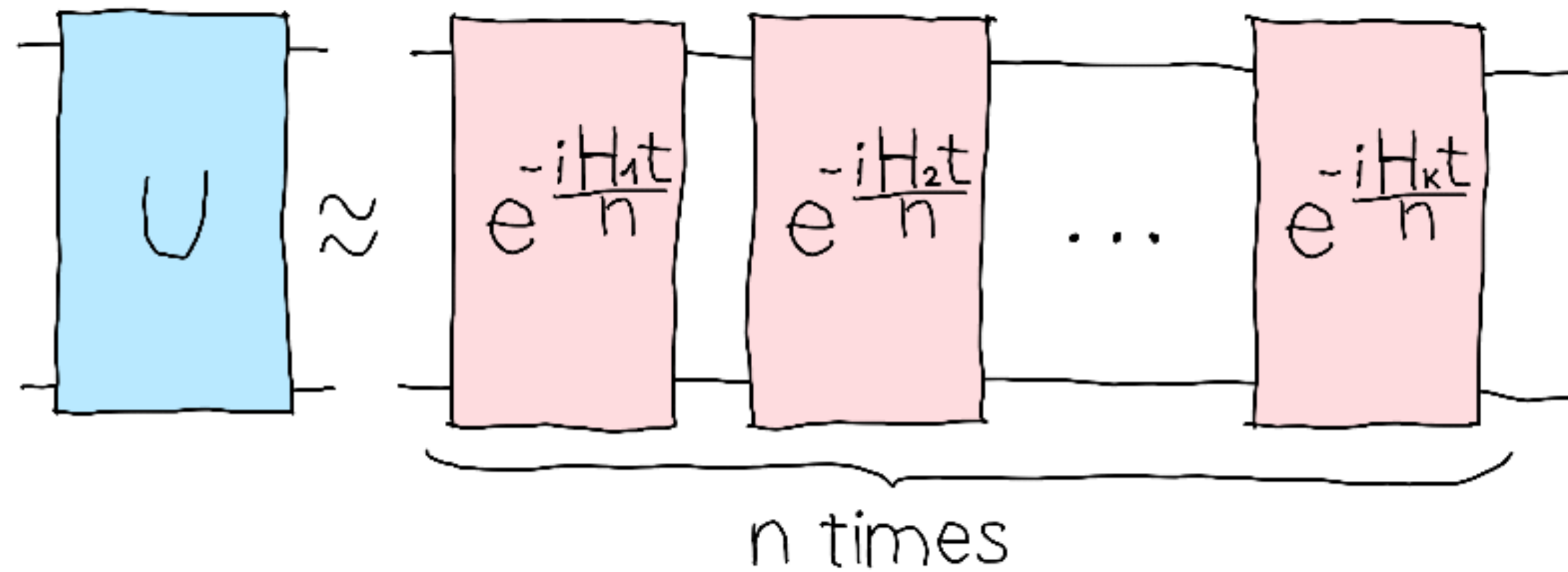
$$\rho_t = e^{-iLt} \rho_0$$

# Simulating coupled quantum-classical dynamics

- Separate Liouvillian into classical and electronic parts, and evolve with small time steps individually.
- No more hybrid loop, but alternating time steps on a quantum computer



# Case closed again?



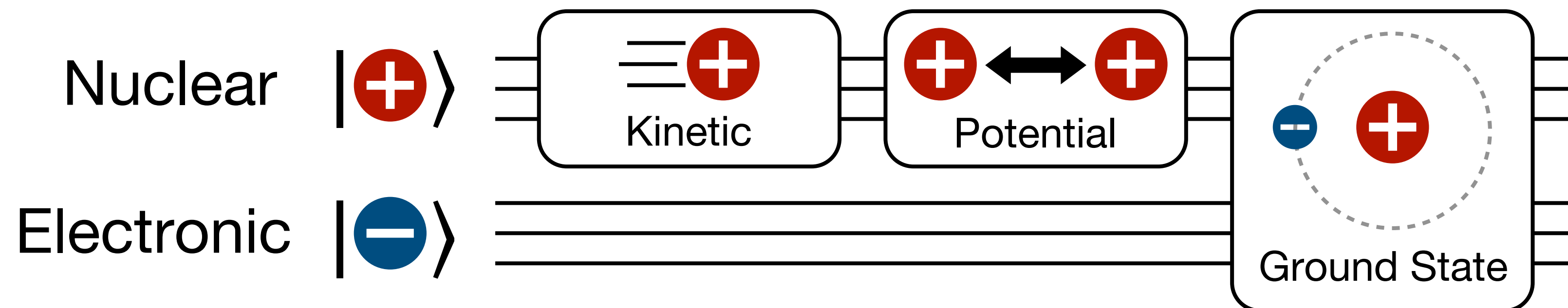
Can we do better than Trotterizing the Liouvillian?



# Directly implementing the Liouvillian

*Or, how to block-encode the Liouvillian*

- Two register sets on quantum computer to implement the Liouvillian:
  - “Classical” nuclear register – Corresponds to classical MD in Liouvillian formalism on exponential sized grid
  - “Quantum” electronic register – Corresponds to ground state energy/forces calculations



$$L = -i \left( \frac{\partial H}{\partial p} \frac{\partial}{\partial x} - \frac{\partial H}{\partial x} \frac{\partial}{\partial p} \right)$$

# Writing down the terms

*Aka the slide with a lot of math*

$$H_{\text{ext}}^{(\text{NVT})} = \sum_{n=1}^N \sum_{j=1}^3 \frac{p'_{n,j}{}^2}{2m_n s^2} + \sum_{k=1}^{N-1} \sum_{n=k+1}^N \frac{Z_n Z_k}{\|x_n - x_k\|} + \frac{p_s^2}{2Q} + N_f k_B T \ln(s) + E_{\text{el}}(\mathbf{x}),$$

$$L_{\text{cl, disc}} := -i \sum_{n=1}^N \sum_{j=1}^3 \left( D_{x_{n,j}} \otimes \frac{\partial H_{\text{cl}}}{\partial p'_{n,j}} - \frac{\partial H_{\text{cl}}}{\partial x_{n,j}} \otimes D_{p'_{n,j}} \right) - i \left( D_s \otimes \frac{\partial H_{\text{cl}}}{\partial p_s} - \frac{\partial H_{\text{cl}}}{\partial s} \otimes D_{p_s} \right), \quad (\text{A.11})$$

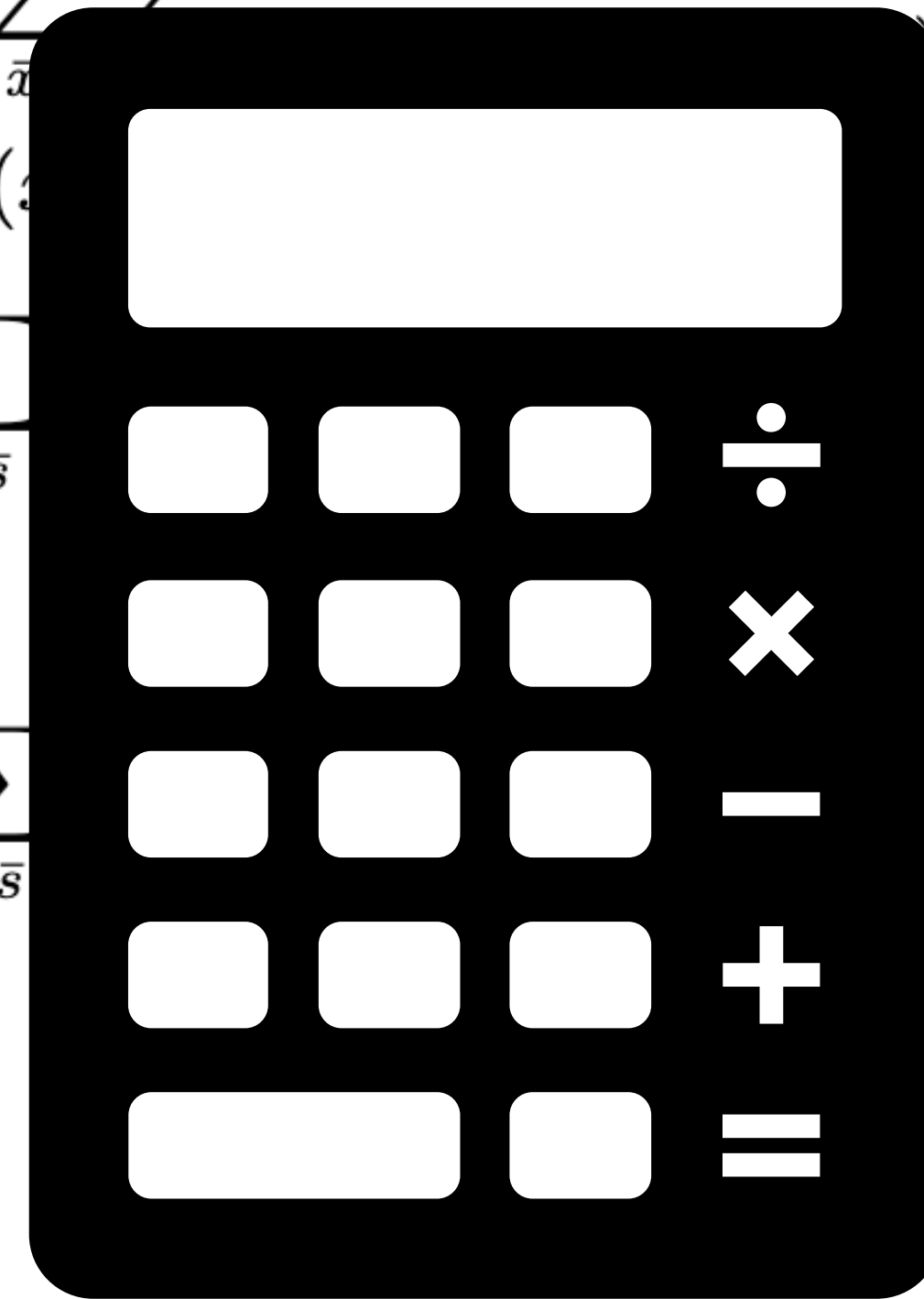
$$L_{\text{el}} = L - L_{\text{cl}} := i \sum_{n=1}^N \sum_{j=1}^3 \frac{\partial E_{\text{el}}(\mathbf{x})}{\partial x_{n,j}} \partial_{p'_{n,j}}$$

$$\frac{\partial H_{\text{cl}}}{\partial x_{n,j}} = \sum_{n' \neq n} \sum_{\bar{x}} \sum_{\bar{x}'} \frac{-Z_n Z_{n'}}{\|x_n - x_{n'}\|^{3/2}} \times (\dots) \langle \bar{x}_{n'} |$$

$$\frac{\partial H_{\text{cl}}}{\partial p'_{n,j}} = \sum_{\bar{p}'_{n,j}} \sum_{\bar{s}} (\dots) | \otimes | \bar{s} \rangle \langle \bar{s} | \quad (\text{A.13})$$

$$\frac{\partial H_{\text{cl}}}{\partial s} = \sum_{\bar{p}_{n,j}} \sum_{\bar{s}} (\dots) | \otimes | \bar{s} \rangle \langle \bar{s} | + \frac{N_f k_B T}{s_{\text{min}}} | \bar{s} \rangle \langle \bar{s} | \quad (\text{A.14})$$

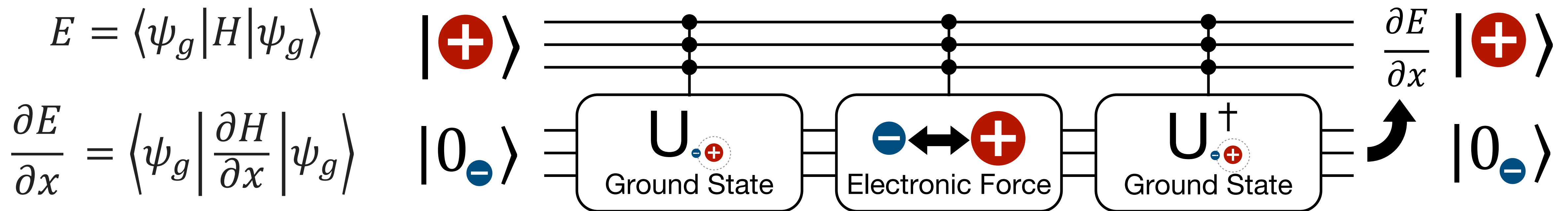
$$\frac{\partial H_{\text{cl}}}{\partial p_s} = \sum_{\bar{p}_s} \frac{p_s}{Q} | \bar{p}_s \rangle \langle \bar{p}_s | \quad (\text{A.15})$$



# Block-encoding energies and forces

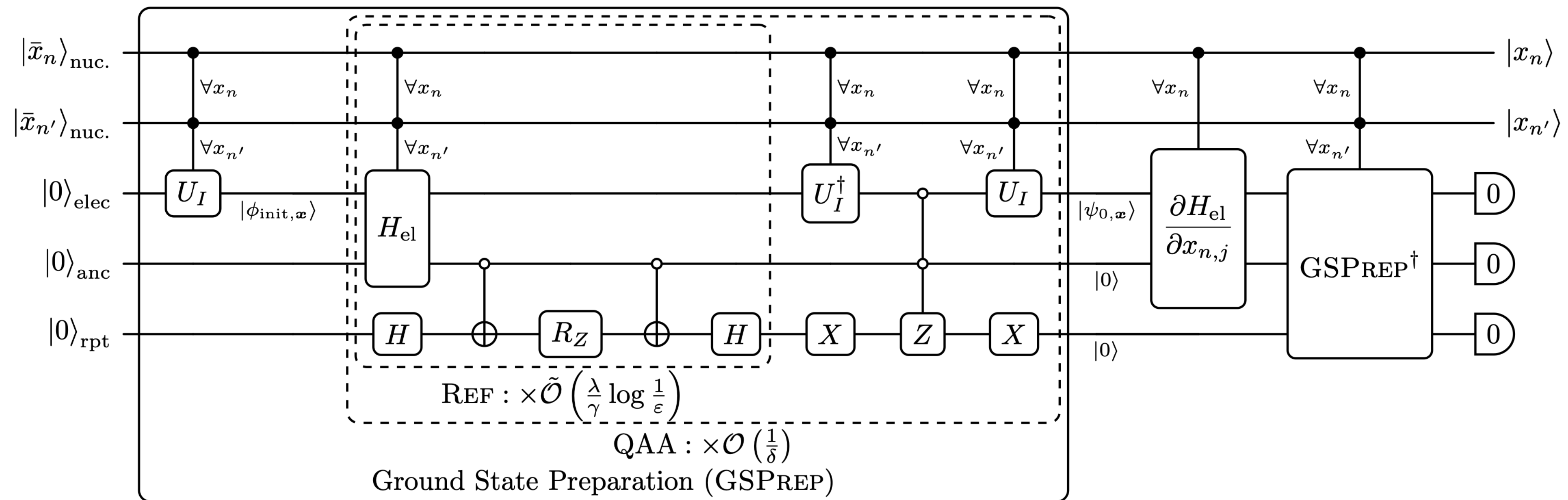
*Obtaining these values without QPE*

- Obtain quantum mechanical electronic forces via kicking back results from electronic register to nuclear register
- Compute ground state energy by preparing ground state, applying the Hamiltonian, then unpreparing the ground state
- Compute electronic forces with Hellmann-Feynman theorem and analytical derivative of first-quantized Hamiltonian



# Implementation of energy derivatives

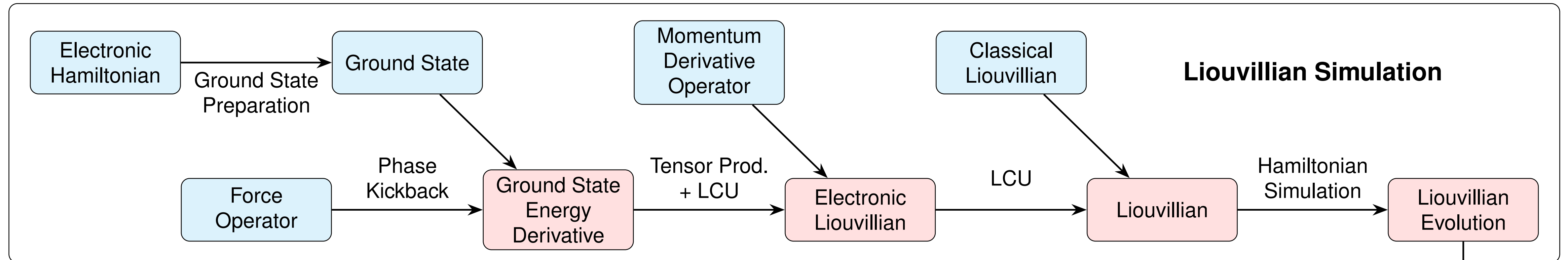
- Block-encoding of first-quantized Hamiltonian derivative found in [OSR+22]
- Ground state preparation algorithm by [LT20]: Eigenstate filtering+AA



# Tying everything together...

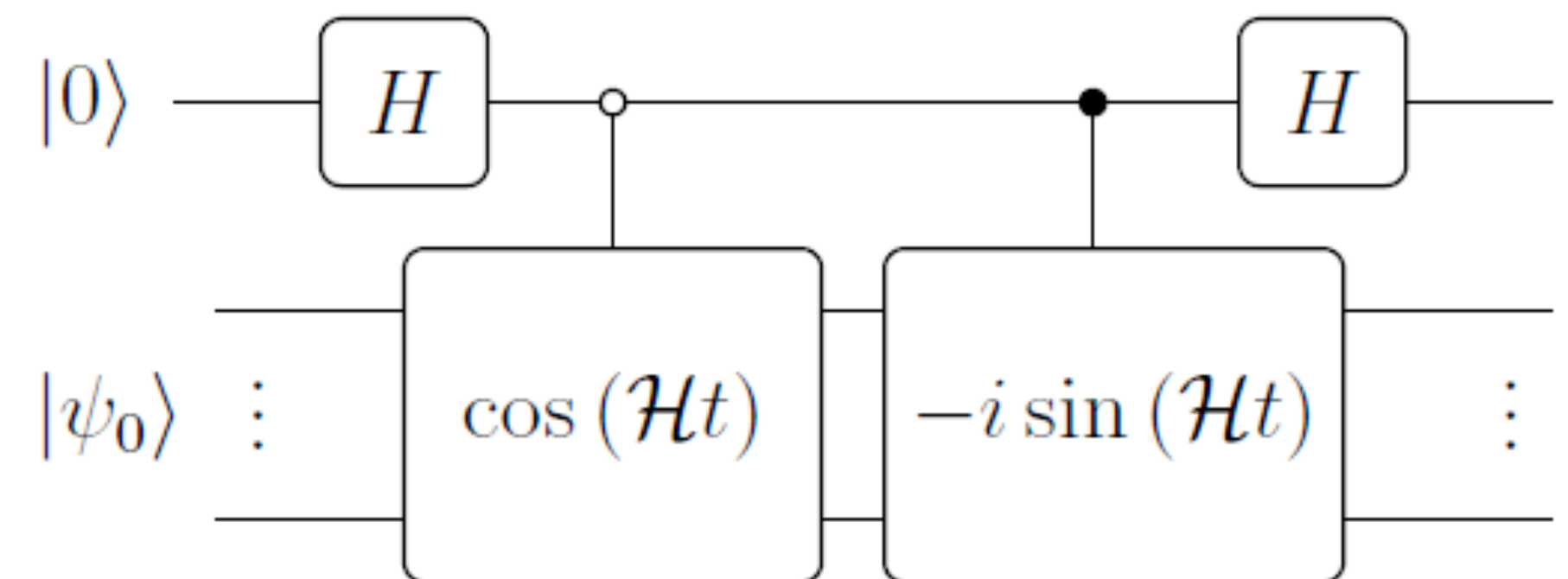
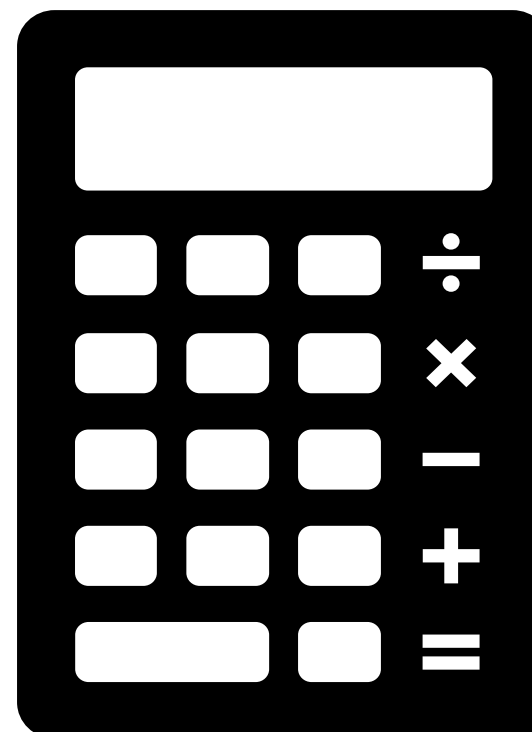
*In Part I...*

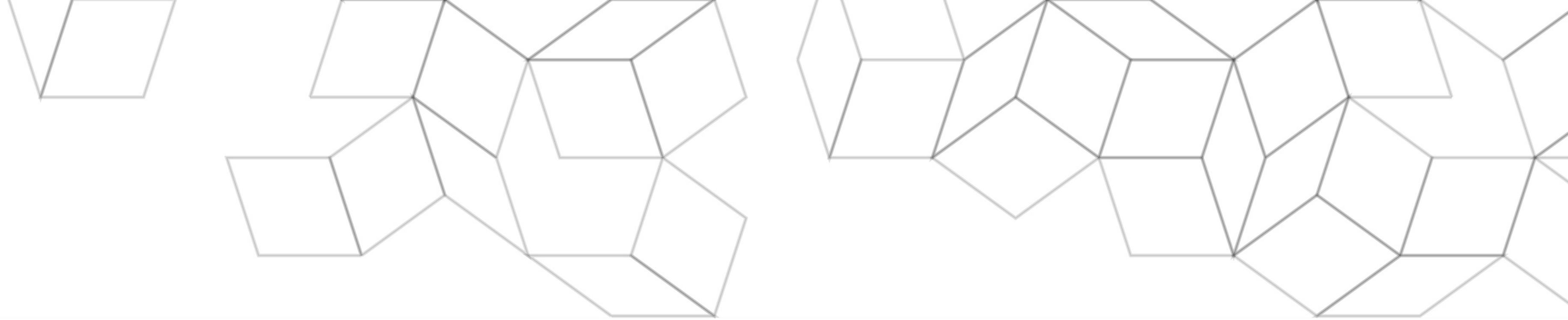
- Liouvillian simulation via QSVT-based Hamiltonian simulation



$$\frac{\partial E}{\partial x} = \left\langle \psi_g \left| \frac{\partial H}{\partial x} \right| \psi_g \right\rangle$$

$$L_{\text{el}} = L - L_{\text{cl}} := i \sum_{n=1}^N \sum_{j=1}^3 \frac{\partial E_{\text{el}}(\mathbf{x})}{\partial x_{n,j}} \partial_{p'_{n,j}}$$





# **Part II:**

# **Alchemical Free Energy Calculation**

# Thermodynamic integration

*The alchemical way to calculate free energy*

$$F = -k_B T \ln Z$$

$$\Delta F_{A \rightarrow B} = \int_0^1 \frac{\partial F(\Lambda)}{\partial \Lambda} d\Lambda$$

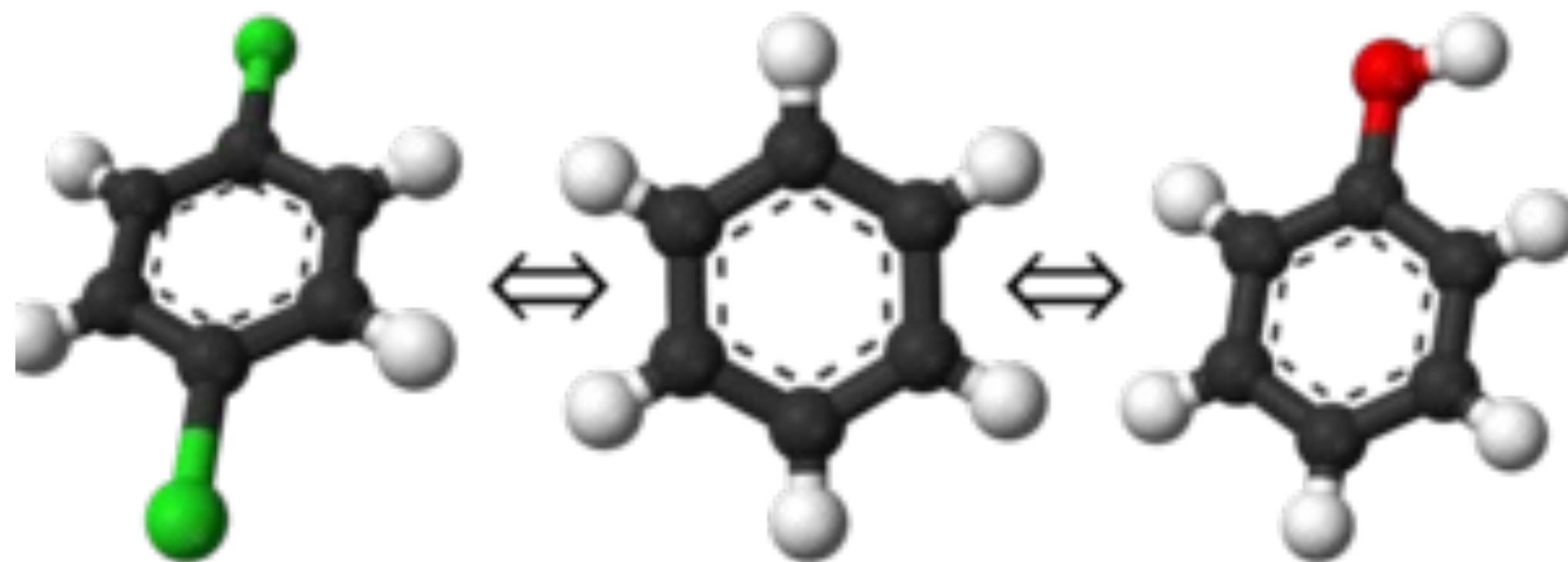
- Absolute free energy difficult to compute
  - Either compute the entropy or partition function

$$= -k_B T \int_0^1 \frac{1}{Z} \frac{\partial Z}{\partial \Lambda} d\Lambda$$

$$= \int_0^1 \frac{1}{Z} \sum_s \frac{\partial E_s(\Lambda)}{\partial \Lambda} e^{-E_s(\Lambda)/k_B T} d\Lambda$$

$$= \int_0^1 \left\langle \frac{\partial E(\Lambda)}{\partial \Lambda} \right\rangle_{\Lambda} d\Lambda,$$

$$= \int_0^1 \langle E(B) - E(A) \rangle_{\Lambda} d\Lambda.$$

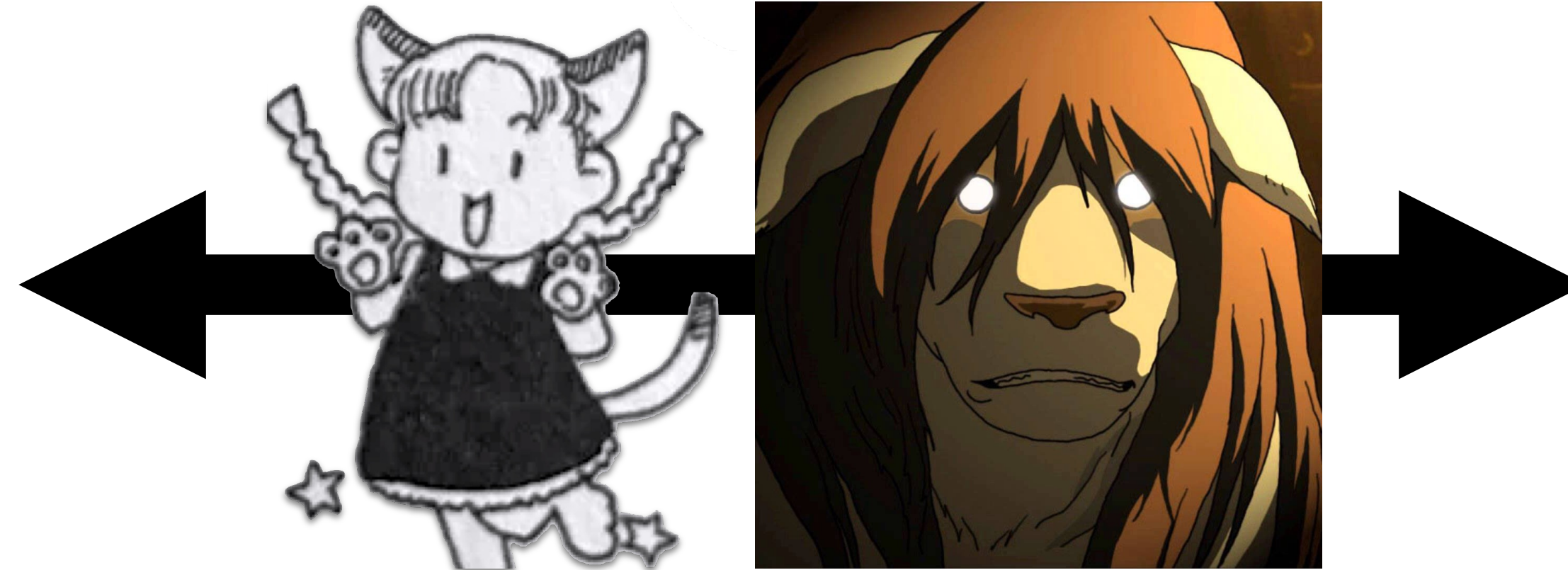


# Thermodynamic integration: A sketch

System A



System B



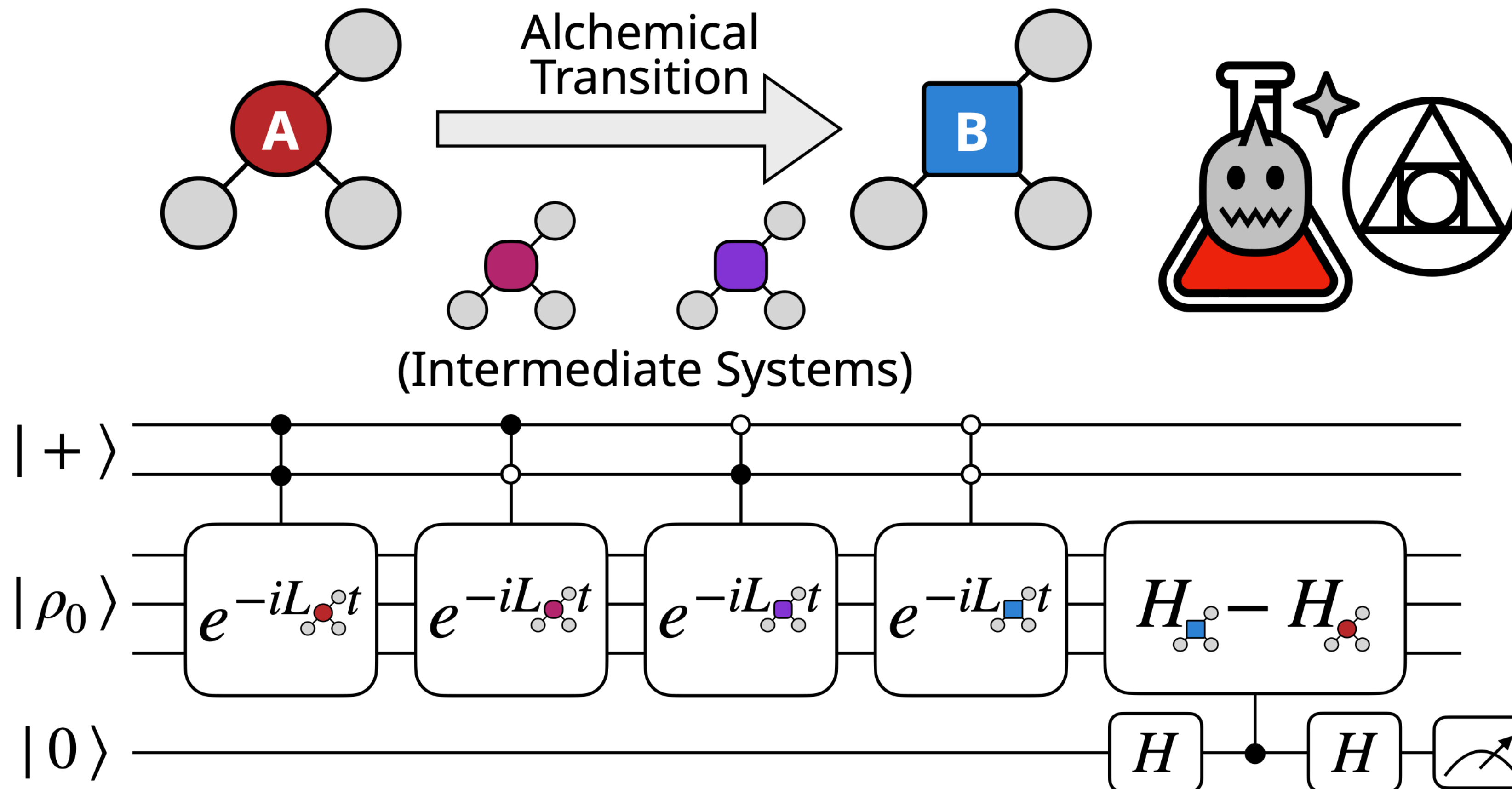
(Non-physical) intermediate state

Generate the thermal distributions under each (intermediate) system

Calculate internal energy difference for each (intermediate) system

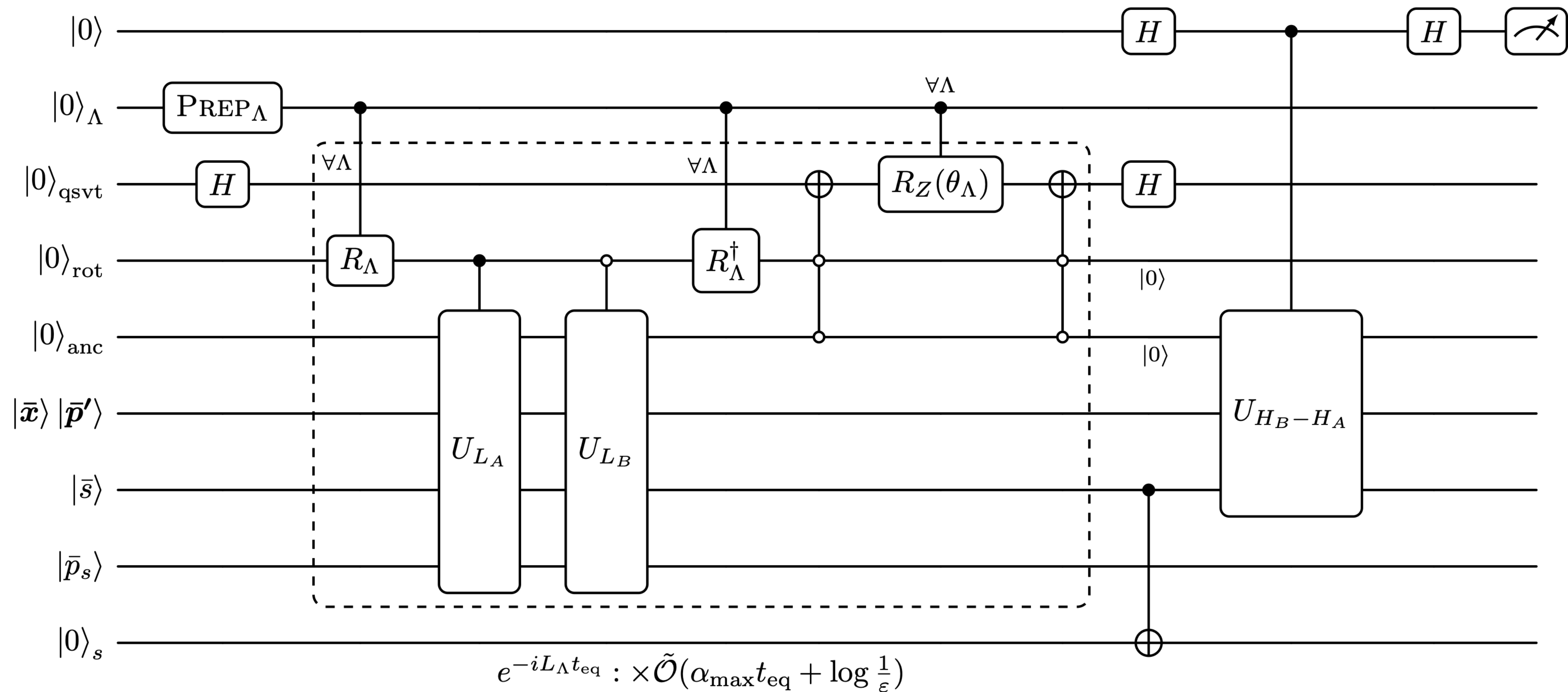
# Calculating free-energy differences

*Thermodynamic integration on a quantum computer*

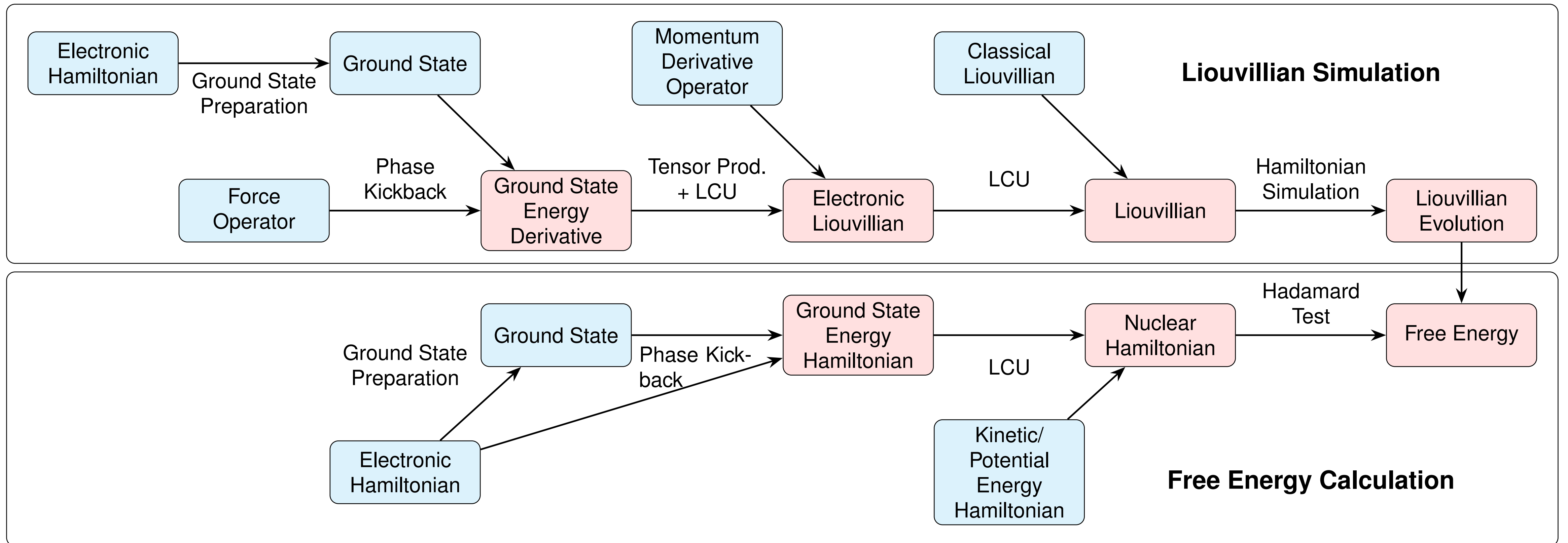


# Controlled Liouvillian simulations

*How the algorithm doesn't scale linear to the interpolation points*



# Algorithm breakdown



# Performance comparison

*With other quantum algorithms and advantages over classical*

- Super-polynomial improvements in precision over iterating between MD time step and electronic structure calculations (with both classical and quantum computing methods)
- Super-polynomial improvements over previous Trotterized Liouvillian approach
- Bypass entropy estimation via relative free energy calculations (no sampling over exponential phase-space)

	MD simulation	Free energy calculation
QPE + Euler Int. [1]	$\tilde{O}(\varepsilon^{-1})$	—
Liouville + Trott. [2]	$\tilde{O}(\varepsilon^{-o(1)})$	$\tilde{O}(\eta^{1+o(1)}\varepsilon^{-1.5} + \varepsilon^{-2})$
Our work	$O(\log^3 \varepsilon^{-1})$	$\tilde{O}(\varepsilon^{-1})$